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NASA TECHNICAL MEMORANDUM

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RADIATIVE AND FREE CONVECTIVE HEAT TRANSFER
FROM A CONTAINERLESS SPHERE

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NASA

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TECHNICAL MEMORANDUM

RADIATIVE AND FREE CONVECTIVE HEAT TRANSFER FROM A CONTAINERLESS SPHERE

INTRODUCTION

Containerless processing of copper in an electromagnetic levitation device uses the natural heating of the sample (generated by the induced eddy currents) and cooling by a helium-argon gas mixture. Such a system has been in operation in the Space Sciences Laboratory, Marshall Space Flight Center, and significant undercooling during solidification has been noted. Temperatures on the order of 2000 K were known to be reached when approximately 20 torr of argon were present in the levitation chamber, and small amounts of helium (approximately 100 torr pressure) were known to reduce the copper sphere's temperature by more than 1000 K. As the coolant gases were admitted to the evacuated chamber through a valve at a large relative distance to the sphere, heat transfer could be considered to be accomplished by radiation and free convection.

In support of this experimentation, a relatively simple mathematical model for the prediction of heat transfer from the sphere was derived. Primary conditions to be considered were the nature of heat loss, the high temperatures and low pressures involved, and the restrictions on direct experimental measurements of various parameters (sphere diameter, emissivity and temperature, and gas temperature and pressure were the only measurable quantities). Parameters to be varied in the calculation of temperature were power input to the levitation coil, emissivity and diameter of the sphere, and gas pressures.

DETERMINATION OF HEAT TRANSFER FORMULA

Power input (PI) to the sphere is equal to total heat loss through radiation and free convection:

$$PI = Q_{\text{rad}} + Q_{\text{conv}}$$

The heat transfer due to radiation (Q_{rad}) is

$$Q_{\text{rad}} = \epsilon \sigma A (T_1^4 - T_o^4)$$

where

$$\sigma = \text{Stefan-Boltzman constant} = 0.56697 \times 10^{-8} \text{ W/m}^2\text{-K}^4$$

$$\epsilon = \text{emissivity of sphere}$$

$$A = \text{surface area of sphere, cm}^2$$

$$T_1 = \text{sphere temperature, K}$$

$$T_o = \text{ambient gas temperature, K}$$

The heat transfer due to free convection (Q_{conv}) is

$$Q_{\text{conv}} = h_c A (T_1 - T_o)$$

where h_c = heat transfer ratio for free convection [1] or, putting h_c in terms of the Nusselt number, Nu,

$$Q_c = A (T_1 - T_o) \frac{k}{D} \text{Nu}$$

where k = thermal conductivity of coolant gas and D = sphere diameter.

Yuge [2] empirically determined the Nusselt number for a sphere in air with heat loss due to free convection to be

$$\text{Nu} = 2 + 0.39 \text{Gr}^{1/4} \quad \text{for} \quad 1 < \text{Gr} < 10^5$$

where Gr = Grashof number. For Grashof numbers much less than one, Yuge recommended Mahony's equation:

$$\text{Nu} = 2 + O (\text{Gr}^{1/2})$$

where O = order of magnitude. By fitting a line to the lower portion of Yuge's data, O was found to equal 0.25. In his work, Mahony dealt

mainly with Grashof numbers between 10^{-3} and 10^{-9} [3]. Because the Grashof numbers encountered fell between 10^{-3} and one, both formulas were tried. No significant differences in results were found. Yuge's equation was used because of the greater documentation and certainty in the Grashof coefficient.

The Grashof number was calculated in the same manner as Yuge:

$$Gr = \frac{gD^3\Delta T}{T_o \nu_f^2}$$

where

g = acceleration due to gravity

D = sphere diameter

ΔT = sphere temperature T_1 - gas temperature T_o

T_o = gas temperature; $1/T_o$ = heat transfer coefficient (β)

ν_f = kinematic viscosity at film temperature $\frac{T_1 + T_o}{2}$

$$= \frac{\rho_f}{\mu_f}$$

where

ρ_f = gas density at film temperature

μ_f = dynamic viscosity at film temperature.

PROPERTIES OF COOLANT GASES

Because no direct measurement of gas properties (with the exceptions of pressure and temperature) was possible, methods for estimating the thermal conductivity, dynamic viscosity, and density were either derived or adapted. All properties were found to be highly temperature dependent, and the temperature difference $T_1 - T_o$ between the sphere

and the surrounding gas was very large. Therefore, all properties were evaluated at the film temperature $T_f = (T_1 - T_o)/2$, as recommended by Yuge.

In the low pressure range found in the experiments, thermal conductivity k is not pressure dependent [4]. A temperature dependency of $T_f^{3/4}$ was determined empirically from data given in McAdams [1]:

$$k = k_o (0.0173) \left(\frac{T_f}{273} \right)^{0.75} \quad \text{W/cm K}$$

where

k_o = thermal conductivity of specific gas at 273 K in Btu/hr·ft-°F

0.0173 = conversion factor from Btu/hr·ft-°F to W/cm K.

This formula was weighted to provide more precise correlation with McAdams' data for higher temperatures and helium, the major component gas.

Since mixtures of helium and argon do not involve rotational and vibrational degrees of freedom, Brokaw's simple empirical method for calculation of thermal conductivity [5] could be used. Brokaw utilized the fact that thermal conductivities for mixtures fall between the values found by simple mixing of the component conductivities and inverse mixing of the components:

$$x_1 k_1 + x_2 k_2 \geq k_m \geq \frac{1}{\frac{x_1}{k_1} + \frac{x_2}{k_2}}$$

where x_1, x_2 = mole fractions of component gases. Brokaw represented the mixture conductivity as a combination of simple and inverse mixing:

$$k_m = q k_{SM} + (1 - q) k_{RM}$$

where

$$k_{SM} = x_1 k_1 + x_2 k_2$$

and

$$\frac{1}{k_{RM}} = \frac{x_1}{k_1} + \frac{x_2}{k_2} .$$

Brokaw then fit this curve to experimental data of numerous binary gas mixtures, including helium-argon, for various molar fractions to determine the value of q . The value of q was found to vary with the ratios of the gases [5]:

mole fraction	0	.1	.2	.3	.4	.5	.6	.7	.8	.9	.95	1.0
q	.32	.34	.37	.39	.42	.46	.50	.55	.61	.69	.74	.8

Brokaw's method was found to have approximately the same accuracy as other methods studied and to be much simpler in form. Therefore, Brokaw's method for gas mixtures was used in combination with the author's formula for individual gases to estimate thermal conductivities of the coolant gases.

Because the dynamic viscosities and the densities of the gases were needed in the calculation of the Grashof number, methods to calculate these properties were also needed. For the component gases, dynamic viscosity μ was taken to be

$$\mu = \mu_o \sqrt{\frac{T_f}{373}} \text{ Poise}$$

where μ_o = dynamic viscosity of specific gas at 373 K (P), again by fitting a curve to McAdams' data [1]. A pressure dependency was not incorporated due to the high temperature and low pressure values to be used.

Brokaw's method for calculation of viscosity of gas mixtures was used, as recommended by Reid, Prausnitz, and Sherwood. Brokaw's method is based on Sutherland's approximation

$$\mu_m = \frac{\sum_{i=1}^n \frac{x_i \mu_i}{\sum_{j=1}^n x_j \phi_{ij}}}{\sum_{j=1}^n x_j \phi_{ij}} ,$$

using

$$\phi_{ij} = \left(\frac{\mu_i}{\mu_j} \right)^{1/2} S_{ij} A_{ij}$$

where

S_{ij} = Sutherland constant

A_{ij} = a complicated function of molecular weights of the component gases.

A_{ij} can be taken from a chart [6] or a graph [4] and, for mixtures of nonpolar gases, S_{ij} can be set equal to one. Thus, for mixtures containing nonpolar gases,

$$\mu_m = \frac{\sum_{i=1}^n \frac{X_i \sqrt{\mu_i}}{\sum_{j=1, j \neq i}^n \frac{A_{ij}}{\sqrt{\mu_i}} X_j}}{\sum_{i=1}^n \frac{X_i}{\sqrt{\mu_i}}}$$

This mixture formula has its best results for mixtures of inert gases; overall errors for nonpolar gas mixtures ranged from 0.6 percent to 2.5 percent error [6]. Brokaw's method can also be applied at any temperature, a vital feature for this project.

The densities of the gases were calculated by the formula

$$\rho = \left(\frac{M}{22.4 \times 10^3} \right) \left(\frac{273}{T_f} \right) \left(\frac{P}{760} \right) \text{ g/cm}^3$$

where T_f is in K, P is in torr, and M = molecular weight of gas. The method of simple mixing was used to calculate the mixture density:

$$\rho_m = X_1 \rho_1 + X_2 \rho_2$$

THE COMPUTER PROGRAM

A FORTRAN program was designed to simplify calculations and extend the range of applicability to experimentation. Appendix A presents the program nomenclature, including a cross-reference to the nomenclature used in the text of this report; Appendix B presents a program listing.

In the program, nested DO loops are used to allow for variances of sphere emissivity EP, sphere diameter D, and power input PIN (*Set Experimental Parameters* section) in addition to argon pressure P1. Two separate DO loops are provided to change gas pressure (*Increase Pressure* and *Decrease Pressure* sections), simulating the experimental input and pump down of gas. The *Increase Pressure* loop uses the incremented mole fraction to calculate helium pressure P2 for constant P1 while the *Decrease Pressure* loop steadily decreases the total pressure PRTTL, holding the molar ratio of the gas mixture constant. Each pressure loop includes an implied loop within which the gas properties and sphere temperature are calculated (*Iterate to Satisfy Temperature Equation* sections). The gas properties are determined by solving the equations presented earlier in this report; the temperature is found using an iterative technique wherein the basic heat transfer is solved for sphere temperature T by using an estimated sphere temperature T_1 . T_1 is initialized at RADT, the sphere temperature should only radiative heat transfer be present. The T_1 's are revised as the implied loop is repeated until the difference $|T - T_1|$ is less than an input tolerance (in this case, 10 K). An iteration check (FLAG, FLAGM) is also provided to check for bad convergence.

RESULTS AND CONCLUSIONS

As pressure is increased, convection takes on an increasingly important role in cooling the sphere. Once the helium enters, convection becomes dominant. Varying the amount of argon initially present does not significantly change the sphere temperature (Figs. 1 and 2). The percentage of helium present in the mixture, and not the total amount of helium, determines the sphere temperature: 90 torr of helium added to 10 torr of argon changes the sphere temperature from 1577 K to 637 K, while 360 torr helium added to 40 torr argon effects a change from 1519 K to 618 K. This phenomenon is due to the fact that the Grashof number is the only pressure-dependent term in the heat transfer equation. As expected, emissivity produced the majority of its effect at

low pressure, where Rad T is proportional to emissivity (Fig. 3). Sphere temperature is directly proportional to power input (Fig. 4) and inversely proportional to sphere diameter (Fig. 5). Once convective cooling dominates, an increase of 1.5 W in power is comparable to a decrease of 0.135 cm in diameter. During radiation's dominance, the power input change will not be as effective as the change in diameter.

Small fluctuations in the value of k_{mix} will effect great changes in the values of Comb T, while Rad T will remain constant. It is believed that the major source of error in this program is in the value of k_{mix} . As temperatures increase, the estimation for individual thermal conductivities will tend to be high. This deviation will decrease the apparent pressure effect on sphere temperature Comb T.

Correlation for the *Increase Pressure* portion of the program with experimental results was found to be good: the variance between sphere temperature at 20 torr argon and 200 torr total was 2.7 percent. The actual temperature values were approximately 5 percent higher than expected; however, the error in calculation of thermal conductivity would predict a higher power input (and thus temperature) for the same cooling effect on the sphere.

Results for the *Decrease Pressure* portion of the program were discouraging. The increase in temperature is only approximately one-third that found experimentally. This is believed to be due to the dominance of the conductivity term, which depends upon the ratio of the gases rather than the total amount of gas present (Fig. 6). The small temperature change shown illustrates the temperature dependence on the Grashof number.

NOTE TO FIGURES:

Because total heat loss did not vary significantly with the total amounts of gas present but rather with the percentage of helium present (Figs. 1,2), Figures 2 through 5 are plotted as temperature versus percentage helium. All data (except where noted) were taken from runs with 20 torr argon initial pressure. The following is an approximation of results for runs beginning with 10 or 40 torr argon:

<u>He, Percent</u>	<u>Total Pressure</u>		
	<u>10 torr Ar</u>	<u>20 torr Ar</u>	<u>40 torr Ar</u>
0	10.0	20.0	40.0
10	11.1	22.2	44.4
20	12.5	25.0	50.0
30	14.3	28.6	57.1
40	16.7	33.3	66.7
50	20.0	40.0	80.0
60	25.0	50.0	100.0
70	33.3	66.7	133.3
80	50.0	100.0	200.0
90	100.0	200.0	400.0

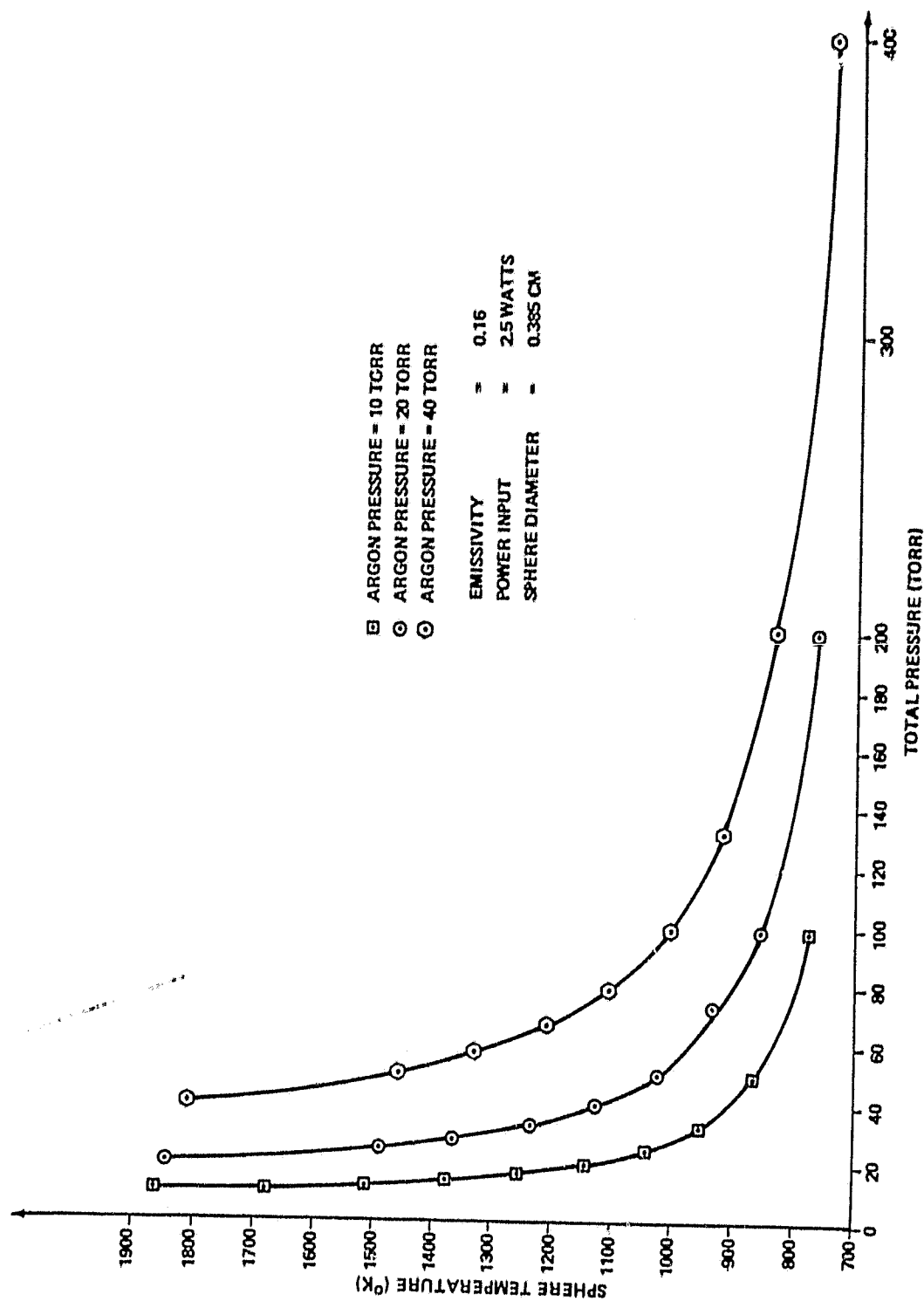


Figure 1. Sample temperature as a function of total pressure for different initial argon pressures.

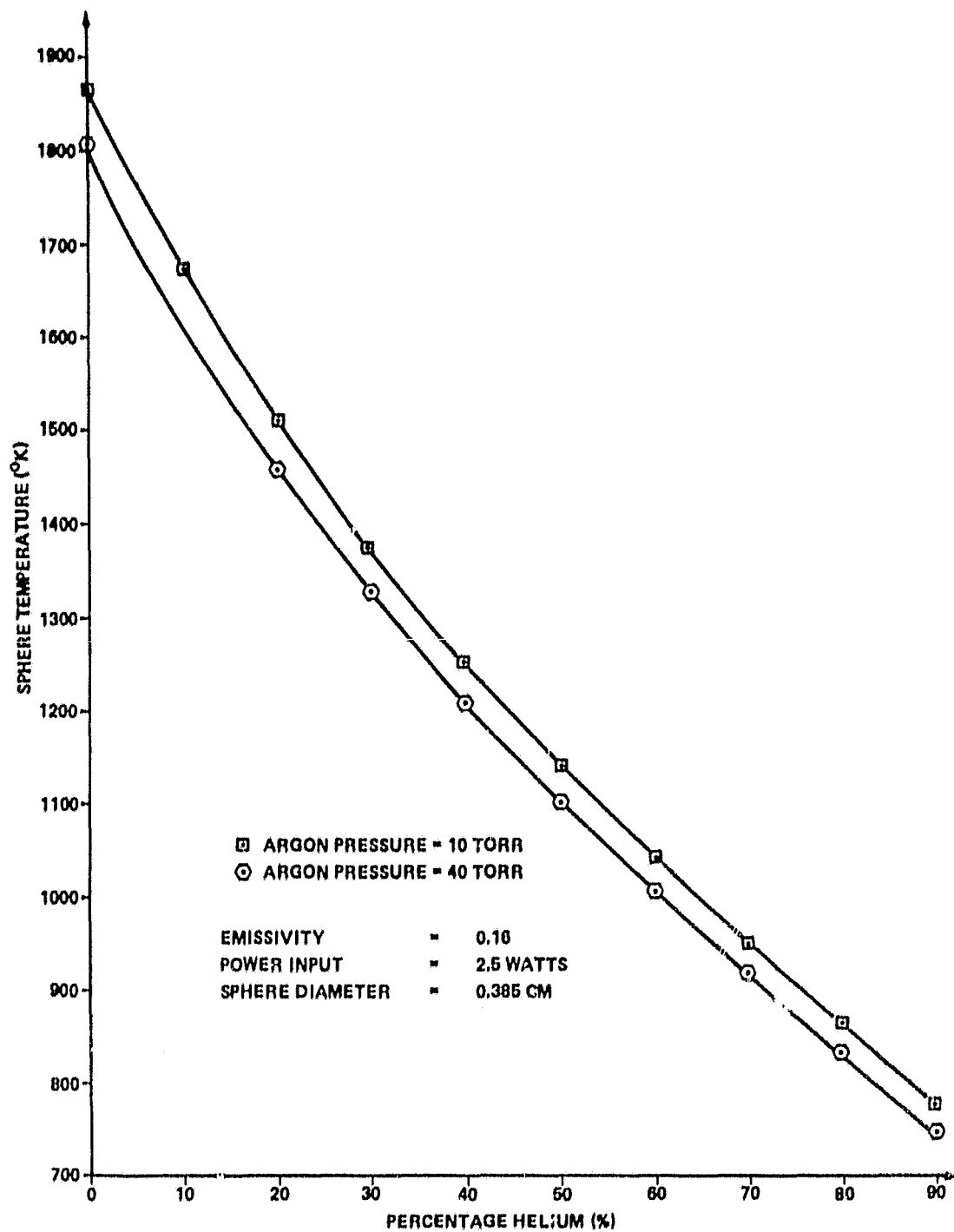


Figure 2. Sample temperature as a function of percentage helium for different initial argon pressures.

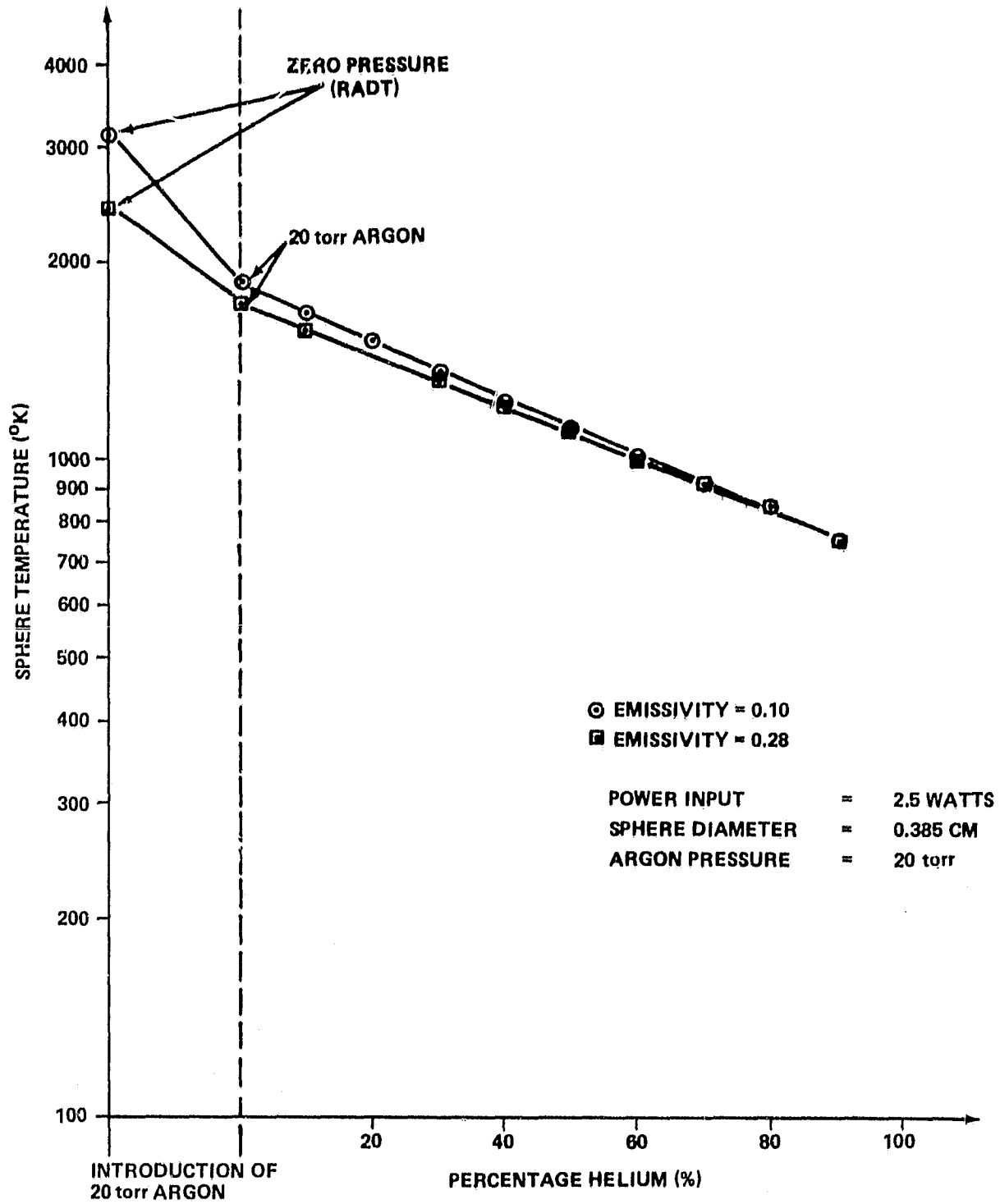


Figure 3. Sample temperature as a function of percentage helium for various emissivities.

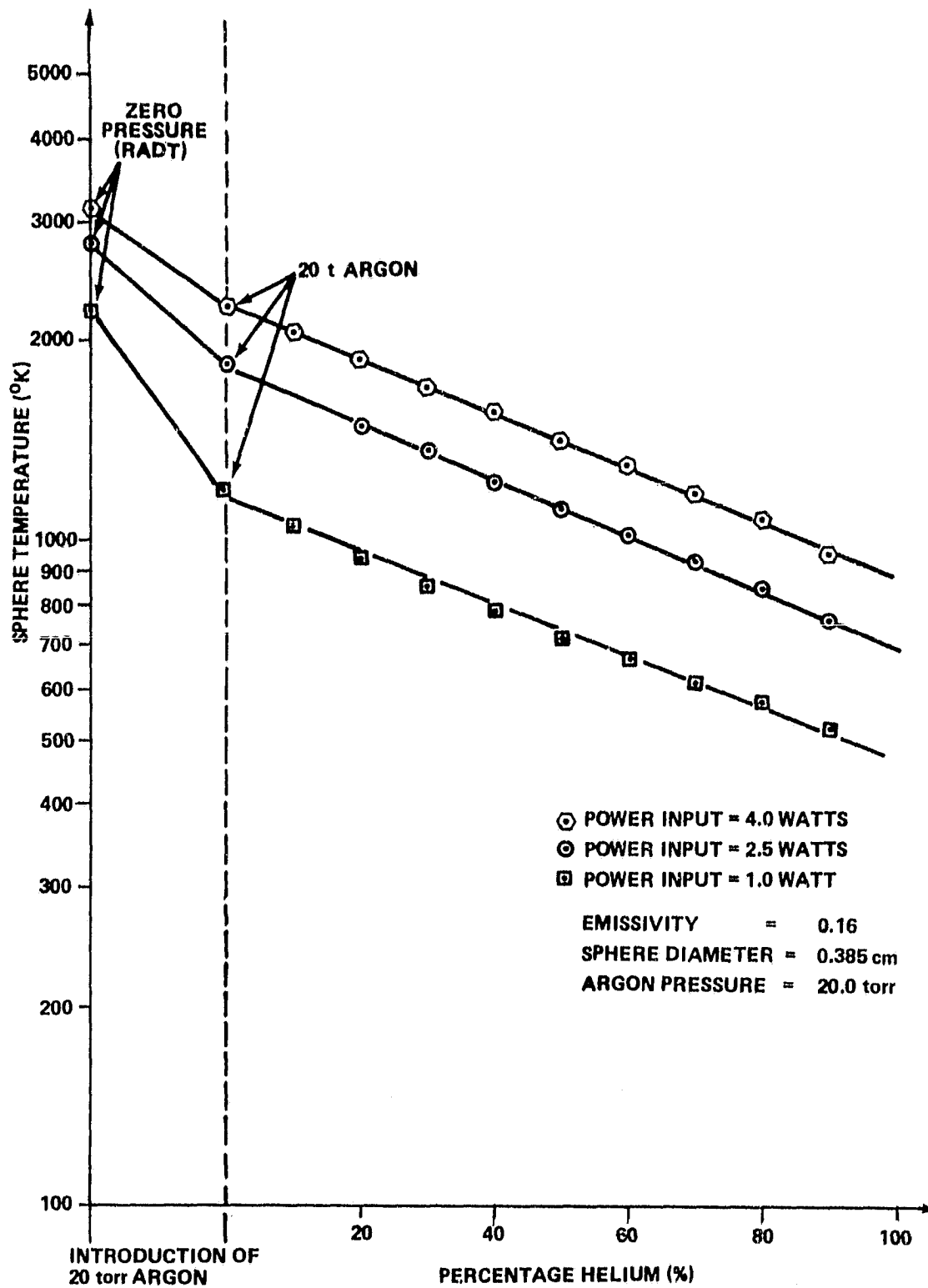


Figure 4. Sample temperature as a function of percentage helium for various power inputs.

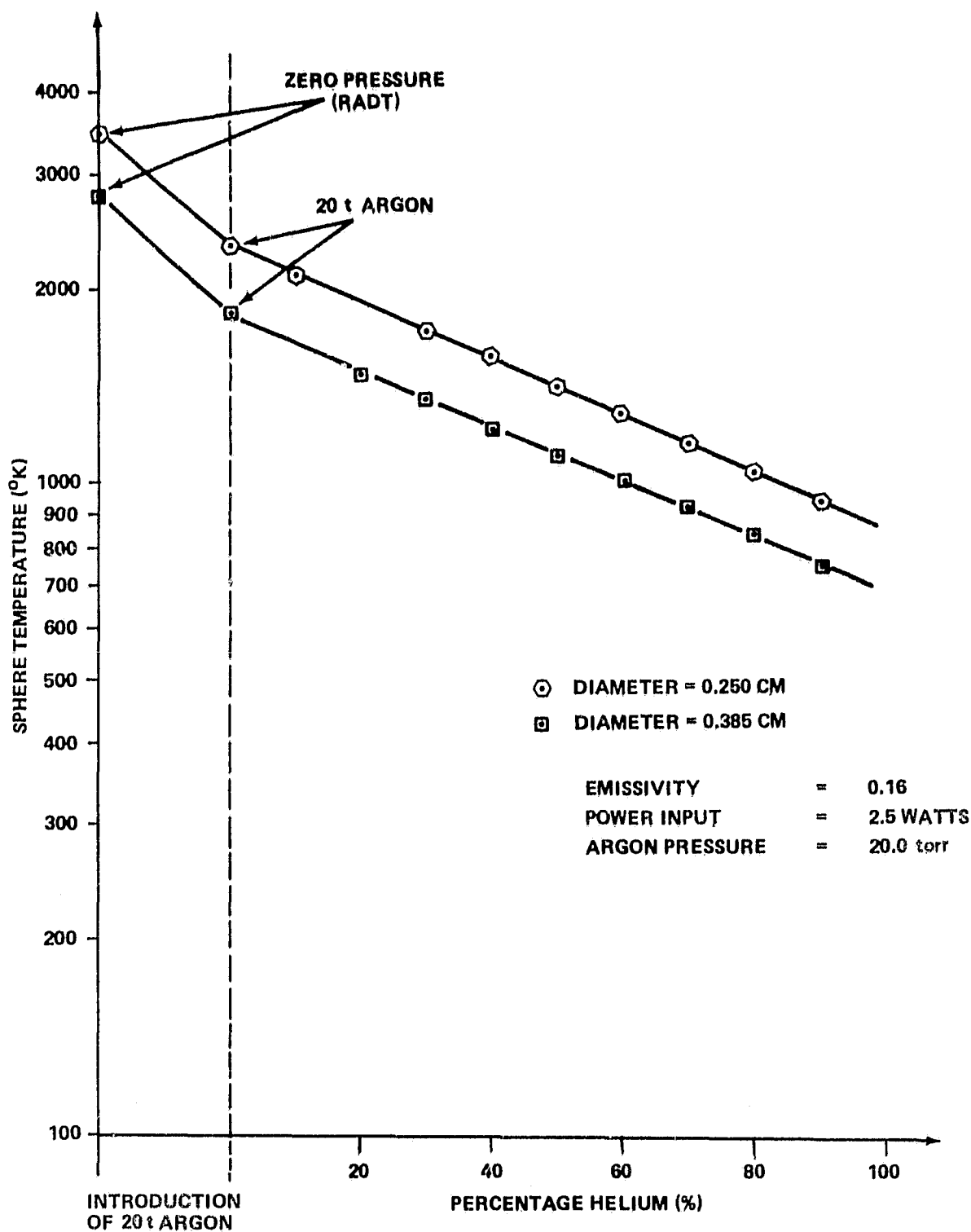


Figure 5. Sample temperature as a function of percentage helium for various sphere diameters.

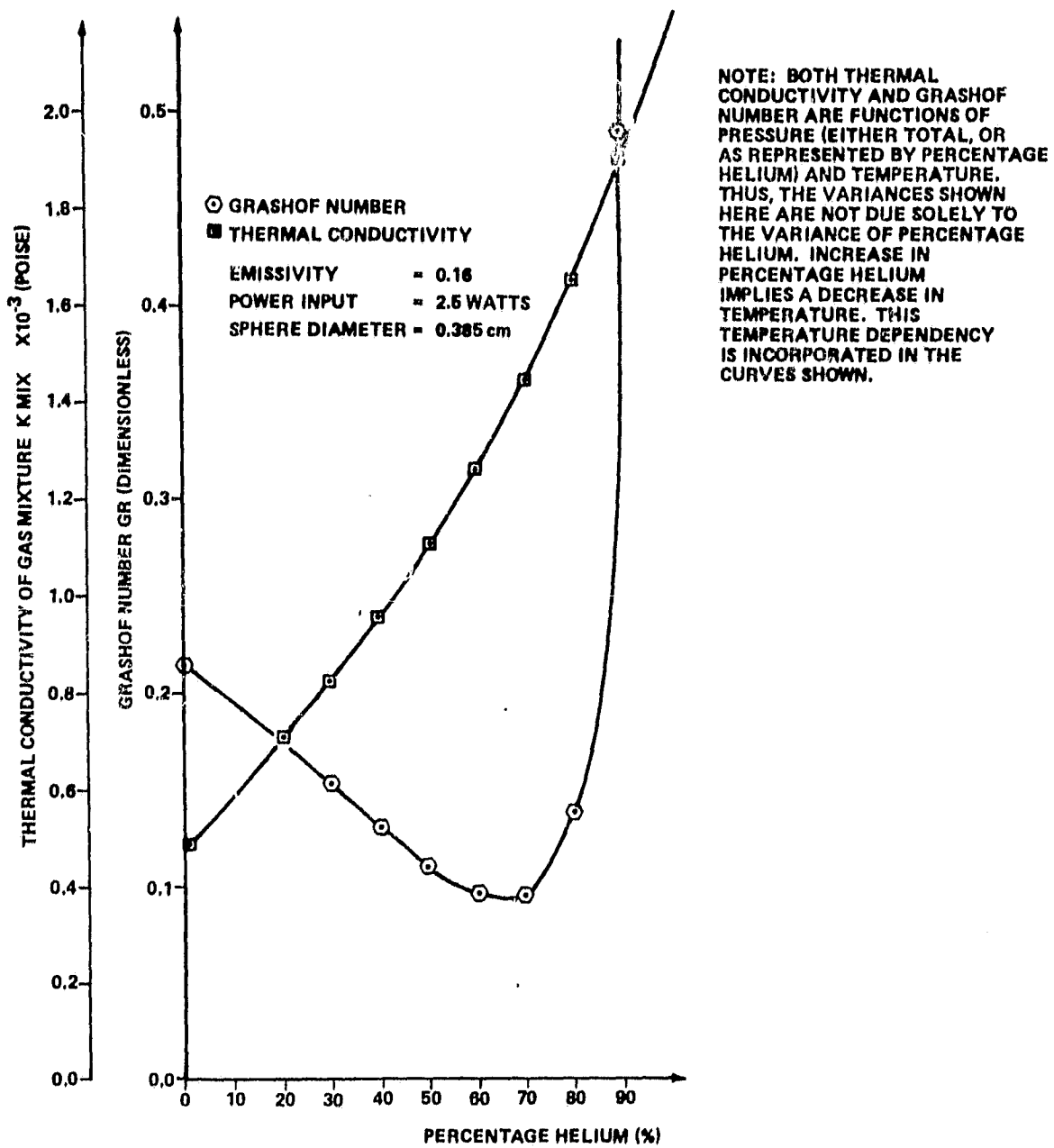


Figure 6. Thermal conductivity and Grashof numbers as functions of percentage helium.

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1. McAdams, W. H.: Heat Transmission. 3rd edition, McGraw-Hill Book Co., Inc., New York, 1954, pp. 165, 457, 468-469.
2. Yuge, T.: Experiments on Heat Transfer from Spheres Including Combined Natural and Forced Convection. Journal of Heat Transfer, vol. 82C, 1960, pp. 214-220.
3. Mahony, J. J.: Heat Transfer at Small Grashof Numbers. Royal Society of London Proceedings, Series A, vol. 238, 1951, pp. 412-423.
4. Reid, Robert C.; Prausnitz, John M.; and Sherwood, Thomas K.: The Properties of Liquids and Gases. 3rd edition, McGraw-Hill Book Co., Inc., New York, 1977, pp. 501, 416-417.
5. Brokaw, Richard S.: Estimating Thermal Conductivities for Nonpolar Gas Mixtures. Industrial and Engineering Chemistry, vol. 47, November 1955, pp. 2398-2400.
6. Brokaw, Richard S.: Viscosity of Gas Mixtures. NASA TN D-4496, 1967, pp. 6-7.

APPENDIX A

NOMENCLATURE

<u>Report</u>	<u>Program</u>	<u>Definition</u>
A_{12}, A_{21}	A12, A21	Dimensionless parameter used in calculation of gas mixture viscosities, function of the molecular weight ratios $\left(\frac{M_1}{M_2}, \frac{M_2}{M_1} \right)$
—	ARGT	The argument of T used in calculation of intermediate values of sphere temperatures $T = \text{ARGT}^{1/4} \text{ (K}^4\text{)}$
—	AVDT	Absolute value of intermediate sphere temperature value difference $ T - T_1 \text{ (K)}$
Comb T	COMBT	Sphere temperature due to combined radiative and free convective heat transfer (K)
D	D	Diameter of sphere (cm)
ϵ	EP	Emissivity of sphere
FLAG	FLAG	Iteration counter for temperature calculation
FLAGM	FLAGM	Maximum number of iterations allowed in calculation of sphere temperature
Gr	GR	Grashof number for gas mixture, dimensionless ratio of buoyant to inertial forces
k_o	K10, K20	Thermal conductivity of gas (1,2) at 273 K (Btu-hr-ft-°F), used in calculation of thermal conductivity of gas (W/cm-K) at film temperature
k_1, k_2	K1, K2	Thermal conductivity of gas (1,2) (W/cm-K)
k_m	KMIX	Thermal conductivity of gas mixture (W/cm-K)
M	M1, M2	Molecular weight of gas (1,2)
μ_o	MU10, MU20	Dynamic viscosity of gas (1,2) at 373 K (P) used in calculation of dynamic viscosity of gas (1,2) at film temperature (P)

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NOMENCLATURE (Concluded)

<u>Report</u>	<u>Program</u>	<u>Definition</u>
μ	MU1,MU2	Dynamic viscosity of gas (1,2) at film temperature (P)
μ_m	MUMIX	Dynamic viscosity of gas mixture at film temperature (P)
P_1, P_2	P1,P2	Amount of pressure of gas (1,2) (torr)
—	PCNT1,PCNT2	Percentage of gas (1,2) in mixture (percent)
PI	PIN	Power input to levitation coil (W)
P_{ttl}	PTTL	Total gas pressure during increase of pressure (torr)
PRTTL	PRTTL	Total gas pressure during pump down (torr)
q	Q(I)	Dimensionless parameter used in calculation of thermal conductivity of gas mixture, function of molar fraction of light gas
Rad T	RADT	Sphere temperature when only radiative cooling effects are considered (K)
ρ_1, ρ_2	RH01,RH02	Density of gas (1,2)
ρ_m	RHOMX	Density of gas mixture
T	T	Intermediate sphere temperature value used in determination of sphere temperature (K)
T_o	TO	Ambient gas temperature (K)
T_1	T1	Intermediate sphere temperature value used in determination of sphere temperature (K)
T_f	—	Film temperature $\frac{T_1 + T_o}{2}$
—	TOL	Tolerance between intermediate and final sphere temperature values (K)
X_1, X_2	X1,X2	Mole fraction of gas (1,2)

APPENDIX B

***** 000007 IS ON CR00002 USING 00024 DLN3 R 0000

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0001 PROGRAM FCONV
0002
0003
0004
0005 C DATE: AUGUST 9, 1979
0006
0007 C PROGRAMMER: KAREN JOHNSON PHONE: 453-4699
0008 C ORGANIZATION: ES-71
0009 C PROJECT: ELECTRO-MAGNETIC LEVITATOR WORK FOR: BILL ORAN
0010
0011 C PROGRAM DESCRIPTION: CALCULATES THE SAMPLE TEMPERATURE DUE
0012 C TO RADIATIVE AND FREE CONVECTIVE HEAT TRANSFER IN DEGREES
0013 C KELVIN WHICH A GIVEN POWER INPUT WILL PRODUCE FOR A
0014 C SPHERICAL MASS OF KNOWN DIAMETER AND EMISSIVITY WHEN
0015 C COOLED BY A BINARY GAS MIXTURE OF KNOWN COMPOSITION
0016
0017 C INPUT: DISC
0018
0019 C OUTPUT: LINE PRINTER
0020
0021
0022
0023 C INPUT DATA
0024
0025
0026 DIMENSION Q(10)
0027 REAL K10,K20,MU10,MU20,M1,M2
0028 DOUBLE PRECISION T1,RADT,T,ARGT,GR,RH01,RH02,RH0X
0029 DOUBLE PRECISION K1,K2,KMIX,MU1,MU2,MUMIX
0030 DATA Q / .32, .34, .37, .39, .42, .46, .50, .55, .61, .69 /
0031 DATA M1/40. /, M2/4. /, MU10/.00026/, MU20/.00022/, T0/300. /
0032 DATA K10/.01/, K20/.082/, A12/.312/, A21/2.18/, TOL/10. /, FLAGH/2000. /
0033 DO 340 L=1,3
0034 EXP1 = FLOAT(L)-1
0035 P1 = 10.*2.**EXP1
0036
0037 C WRITE INITIAL DATA
0038
0039 WRITE(6,20) FLAGH,TOL,T0,A12,A21
0040 20 FORMAT(" THE MAX NO OF ITERATIONS FLAGH = ",F5.0/" THE TOLERANCE T
0041 IOL",14X," = ",F5.1/" THE AMBIENT GAS TEMPERATURE T0 = ",F6.1/" VIS
0042 COSITY MIXTURE PARAMETERS"/" A12 = ",F6.3,6X," A21 = ",F6.3//)
0043 WRITE(6,30) M1,M2,MU10,MU20,K10,K20
0044 30 FORMAT(37X,"1",16X,"2"/35X," ARGON",12X," HELIUM"/" MASS",31X,F3.0,
0045 114X,F3.0/" COEF OF VISCOSITY (MU0)",10X,F7.5,10X,F7.5/" COEF OF THE
0046 RM COND (K0)",11X,F3.3,12X,F5.3//)
0047 WRITE(6,40)
0048 40 FORMAT(" EP = EMISSIVITY OF SPHERE"/" PIN = POWER INPUT TO LE
0049 VITATION COIL (WATTS)"/" D = DIAMETER OF SPHERE (CENTIMETERS)"
0050 12" PTTL = TOTAL GAS PRESSURE (TORRS)"/" PCNT1 = PERCENTAGE OF GAS
0051 1 IN MIXTURE"/" PCNT2 = PERCENTAGE OF GAS 2 IN MIXTURE"/" RADT =
0052 SPHERE TEMP DUE TO RADIATIVE HEAT TRANSFER (P = 0, DEG KELVIN)"/"
0053 COMBT = SPHERE TEMP DUE TO COMBINED RADIATIVE AND CONVECTIVE HEAT
0054 TRANSFER (DEG KELVIN)"/" GR = GRASHOF NUMBER FOR GAS MIXTURE"
0055 12//)
0056
0057
0058
0059 C SET EXPERIMENTAL PARAMETERS

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0061      DO 320 IEP=10.25.6
0062      EP = FLOAT(IEP)/100
0063      DO 310 ID=250.385.135
0064      D = FLOAT(ID)/1000
0065      DO 300 IPIN=10.40.15
0066      PIN = FLOAT(IPIN)/10
0067
0068      I
0069      C INCREASE PRESSURE
0070      I
0071      I
0072      WRITE(6,45)
0073 45    FORMAT(' INCREASE PRESSURE')
0074      WRITE(6,50)
0075 50    FORMAT(3X,'EP',8X,'PIN',9X,'D',10X,'PTTL',8X,'PRES1',7X,'PRES2',
0076      ' 7X,'PCNT2',7X,'RADT',8X,'COMBT',10X,'GR',10X,'KMIX'/)
0077      T1 = (PIN/(1.78E-12*D**2*EP)+T0**4)**.25
0078      RADT = T1
0079      DO 180 IX2=1.10
0080      X2 = (FLOAT(IX2)-1.)/10.
0081      X1 = 1.-X2
0082      I = IX2
0083      P2 = X2*(P1/X1)
0084      FLAG = 0
0085      I
0086      C ITERATE TO SATISFY TEMPERATURE EQUATION
0087      C
0088 60    FLAG = FLAG + 1
0089      RHQ1 = (3.21E-5*M1*P1)/(T1+T0)
0090      RHQ2 = (3.21E-5*M2*P2)/(T1+T0)
0091      RHOMX = X1*RHQ1+X2*RHQ2
0092      MU1 = 3.66E-2*MU10*(T1+T0)**.5
0093      MU2 = 3.66E-2*MU20*(T1+T0)**.5
0094      HUMIX = ((X1*MU1)/(X1+X2*(MU1/MU2))+.5*X2))/((X2*MU2)/(X2+X2*(MU2/MU1))+.5*X1))
0095      K1 = 1.53E-4*K10*((T1+T0)**.75)
0096      K2 = 1.53E-4*K20*((T1+T0)**.75)
0097      KMIX = Q(I)*(X1*K1+X2*K2)+(1.-Q(I))*(1./((X1/K1)+(X2/K2)))
0098      IF (X2-3) 65,65,93
0099 65    ARG1=((PIN-(6.28*KMIX*D*(T1-T0)))-((6.86*KMIX*D**1.75*RHOMX**5*(T1-T0)**1.25)/(T0**25*MUMIX**5)))/(1.78E-12*EP*D**2)
0100      +T0**4
0101      IF ARG1 70,70,90
0102 70    IF FLAG-FLAG 140,140,80
0103 80    T1 = T1 - 100
0104      GO TO 60
0105 90    T = ARG1**25
0106      GO TO 96
0107 96    T=((PIN-(1.78E-12*EP*D**2*(T1**4-T0**4)))-((6.86*KMIX*D**1.75*RHOMX**5*(T1-T0)**1.25)/(T0**25*MUMIX**5)))/(6.28*KMIX*D)
0108      +T0
0109 96    AVDT = ABS(T-T1)
0110      IF (TOL-AVDT) 100,160,160
0111 100   IF FLAG-FLAG 140,140,110
0112 110   T1 = (T+T1)/2
0113      IF (T1-T0) 120,120,130
0114 120   T1 = T0 + 100
0115
0116 130   GO TO 60

```

```

0120 C CALCULATE AND WRITE FINAL DATA
0121 C
0122 140 WRITE (6,150) T,T1,RADT
0123 150 FORMAT(," BAD CONVERGENCE THE LAST TWO VALUES OF T WERE ".F6.0,
0124 " AND ".F6.0," RADT = ".F6.0/)
0125 GO TO 180
0126 160 PTTL = P1 + P2
0127 PCNT2 = X2*100.
0128 GR = (980 + D**3*RHOMX**2*(T1-T0))/(T0*MUMIX**2)
0129 WRITE(6,170) EP, PIN, D, PTTL, P1, P2, PCNT2, RADT, T, GR, KMIX
0130 170 FORMAT(F5.2, F11.1, F11.3, F13.1, 3F12.1, F13.0, F12.0, E15.3, E13.3)
0131 180 CONTINUE
0132 C
0133 C
0134 C DECREASE PRESSURE
0135 C
0136 C
0137 WRITE (6,185)
0138 185 FORMAT(//25X,"*****")
0139 "*****"// " DECREASE PRESSURE"//3X,"EP",8X,"PIN",9X,"D",10X,
0140 "PTTL",8X,"PRES1",7X,"PRES2",7X,"PCNT2",7X,"RADT",8X,"COMBT",10X,
0141 "GR",10X,"KMIX"//)
0142 C
0143 C SET CONSTANTS AND INITIALIZE VARIABLES
0144 C
0145 X1 = 1
0146 X2 = 9
0147 PCNT2 = X2 * 100
0148 PRITL = PTTL
0149 T1 = RADT
0150 DO 200 N=1,13
0151 PRES1 = X1*PRITL
0152 PRES2 = X2*PRITL
0153 FLAG = 0
0154 C
0155 C ITERATE TO SATISFY TEMPERATURE EQUATION
0156 C
0157 190 FLAG = FLAG + 1.
0158 RH01 = (3.21E-5*M1*PRES1)/(T1+T0)
0159 RH02 = (3.21E-5*M2*PRES2)/(T1+T0)
0160 RHOMX = X1*RH01+X2*RH02
0161 MU1 = 3.66E-2*MU10*(T1+T0)**.5
0162 MU2 = 3.66E-2*MU20*(T1+T0)**.5
0163 MUMIX = ((X1*MU1)/(X1+A12*(MU1/MU2)**.5*X2))+((X2*MU2)/(X2+A21*
0164 (MU2/MU1)**.5*X1))
0165 K1 = 1.53E-4*K10*((T1+T0)**.75)
0166 K2 = 1.53E-4*K20*((T1+T0)**.75)
0167 KMIX = Q(I)*(X1*K1+X2*K2)+(1.-Q(I))*(1./((X1/K1)+(X2/K2)))
0168 T = ((PIN-(1.78E-12*EP*D**2*(T1**4-T0**4))-((6.86*KMIX*D**1.75*
0169 RHOMX** 5*(T1-T0)**1.25)/(T0** .25*MUMIX** .5)))/(6.28*KMIX*D))
0170 +10
0171 AYD1 = DABS(T-T1)
0172 IF(T01-AYD1) 200,260,260
0173 200 IF(FLAG-FLAG) 240,240,210
0174 210 T1 = (T+T1)/2
0175 IF(T1-T0) 220,220,230
0176 220 T1 = T0 + 100.
0177 230 GO TO 190
0178 C

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0174 * CALCULATE AND WRITE FINAL DATA
0180 C
0181 240 WRITE(6,250) T,T1
0182 250 FORMAT(2* BAD CONVERGENCE T = *,F6.0,* T1 = *,F6.0/)
0183 GO TO 280
0184 260 TR = (980.40+3.4RHOMX+2*(T1-T0))/(T0+MUMIX+2)
0185 WRITE(6,270) EP,PIH,D,PRYL,PRES1,PRES2,PCNT2,RAOT,T,GR,KMIX
0186 270 FORMAT(F5.2,F11.1,F11.3,F13.1,3F12.1,F13.0,F12.0,E15.3,E13.3)
0187
0188 C CHECK IF PUMP DOWN IS COMPLETED
0189 C
0190 PRTL = PRTL-30.
0191 IF(PRTL LE.0.) GO TO 285
0192 280 CONTINUE
0193 284 WRITE(6,290)
0194 290 FORMAT(25X,*)
0195 *****//)
0196 300 CONTINUE
0197 310 CONTINUE
0198 320 CONTINUE
0199 WRITE(6,330)
0200 330 FORMAT(1H1)
0201 340 CONTINUE
0202 STOP
0203 END
0204 END$

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APPROVAL


RADIATIVE AND FREE CONVECTIVE HEAT TRANSFER FROM A CONTAINERLESS SPHERE

By Karen Johnson

The information in this report has been reviewed for technical content. Review of any information concerning Department of Defense or nuclear energy activities or programs has been made by the MSFC Security Classification Officer. This report, in its entirety, has been determined to be unclassified.



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